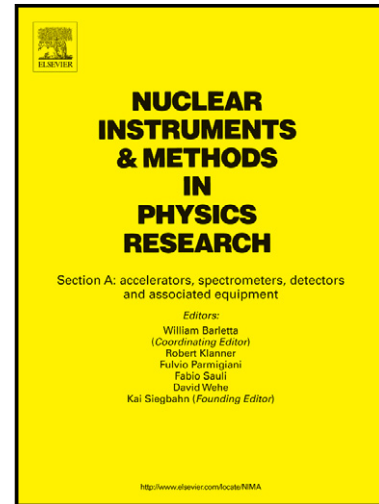


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Beam Studies of the Segmented Resistive WELL: a Potential Thin Sampling Element for Digital Hadron Calorimetry

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Abstract

Thick Gas Electron Multipliers (THGEMs) have the potential of constituting thin, robust sampling elements in Digital Hadron Calorimetry (DHCAL) at future colliders. We report on recent beam studies of new single- and double-THGEM-like structures: the multiplier is a Segmented Resistive WELL (SRWELL) - a single-faced THGEM in contact with a segmented resistive layer inductively coupled to readout pads. Several 10×10 cm² configurations with a total thickness of 5-6 mm (excluding electronics) with 1 cm² pads were investigated with muons and pions. The pads were coupled to a Scalable Readout System APV chip, APV-SRS [22]. Detection efficiencies in the 98% range were recorded with an average pad-multiplicity of ~ 1.1 . The resistive anode resulted in efficient discharge damping, with potential drops of a few volts; the discharge probabilities were $\sim 10^{-7}$ for muons and $\sim 10^{-6}$ for pions, at rates of a few kHz/cm² and for detectors in the double-stage configuration. Further optimization work and research on larger detectors are underway.

Keywords: Micropattern gaseous detectors (MPGD), THGEM, SRWELL, Digital hadron calorimetry (DHCAL), Resistive electrode, SRS, ILC, CLIC

1. Introduction

The Thick Gas Electron Multiplier (THGEM) [1] is a simple and robust electrode suitable for large area detectors, which can be economically produced by industrial Printed Circuit Board (PCB) methods. Its properties and potential applications are reviewed in [2,3]; recent progress can be found in [4-7]. One possible application of THGEM-based detectors is in Digital Hadronic Calorimeters (DHCAL), of the kind proposed for the ILC/CLIC-SiD experiment [8,9]. In this project, the calorimeter design dictates very narrow sampling elements, in the sub-centimeter range, with a lateral pixel size of 1×1 cm². Additional requirements are a high detection efficiency ($>95\%$) and a minimum pad-multiplicity (number of pads activated per particle).

RPCs presently constitute the baseline technology for the SiD DHCAL, with 94% efficiency and a pad-multiplicity of 1.6 [10]; other solutions have been investigated, e.g. MICROMEGAS with 98% efficiency and a multiplicity of 1.1 [11], and Double-GEMs with 95% efficiency and a pad-multiplicity of 1.3 [12] (all results are for muons).

Recently, THGEM-based sampling elements were proposed: they were studied with muons and pions, primarily in single- and double-THGEM configurations with direct charge collection on readout pads, separated from the multiplier by a 2 m induction gap [13]. The potential value of this concept for

DHCAL was demonstrated, leaving room for further optimization, in terms of stability in hadronic beams, efficiency, multiplicity, and overall thickness.

We report here on the results of our latest beam study, conducted at the CERN SPS/H4 RD51 beam-line with 150 GeV/c muons and pions. Further substantial progress was made with a novel THGEM-like configuration, the Segmented Resistive WELL (SRWELL). More details can be found elsewhere [14].

2. Experimental setup and methodology

The SRWELL, first suggested in [13], is shown schematically in Figures 1 and 2; it is a THGEM that is copper-clad on its top side only, whose bottom is closed by a resistive anode. The anode consists of a 0.1 mm thick FR4 sheet patterned with a square grid of narrow copper lines, with the entire area coated with a resistive film (e.g. graphite mixed with epoxy [15]). Avalanche-induced signals are recorded inductively on a pad array located below the FR4 sheet. The grid lines on the resistive anode correspond to the inter-pad boundaries; they serve to prevent charge spreading across neighboring pads by allowing for rapid draining of the avalanche electrons diffusing across the resistive layer. The resistive layer itself (resistivity of ~ 10 - 20 M Ω /square) serves to significantly reduce the energy of occasional discharges. The closed-bottom geometry, also suggested in [16] and similar in its field shape to the WELL [17] and

50 C.A.T. (the French acronym for "Compteur À Trou") [18], re-
 51 duces the total thickness of the detector; it also results in at-
 52 taining a higher gain at a given applied voltage, compared to a
 53 standard THGEM with an induction gap [13]. The SRWELL
 54 has a segmented square hole-pattern with "blind" copper strips
 55 above the grid lines; these prevent more energetic discharges in
 56 the holes located above the metal grid.

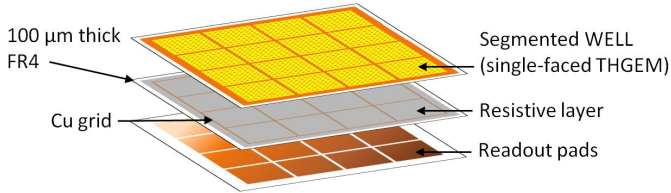


Figure 1: The three layers comprising the SRWELL. Bottom: readout pad array (here 4×4); middle: resistive layer on top of a copper grid (on FR4 sheet); top: segmented single-faced THGEM. The layers are assembled one on top of the other in direct contact (see Fig. 2).

57 Two basic detector configurations were investigated (Fig-
 58 ure 2): one comprising a single-stage SRWELL, and the other,
 59 a double-stage structure with a standard THGEM followed by
 60 an SRWELL. In both cases the electrodes were 10×10 cm² in
 61 size. Based on previous experience with neon-based gas mix-
 62 tures, which allow for high-gain operation at relatively low vol-
 63 tages [4], the detectors were operated in Ne/5%CH₄ at 1 atm,
 64 with a typical flow of a few l/h; a minimally ionizing particle
 65 (MIP) passing through this gas mixture generates, on the aver-
 66 age, ~60 electron-ion pairs per cm along its track [19].

67 In the single-stage detector the SRWELL was either 0.4 m-
 68 m or 0.8 mm thick, with corresponding drift gaps of 5.5 m-
 69 m and 5 mm, respectively. In the double-stage configuration,
 70 both the THGEM and the SRWELL were 0.4 mm thick; the
 71 transfer gap between them was 1.5 mm wide and the drift gap
 72 was 2.5 mm, 3 mm, or 4 mm in width. The total thickness
 73 of the detector from the resistive anode to the drift electrode
 74 was thus between 4.8 and 6.3 mm. The THGEM and the SR-
 75 WELL electrodes were manufactured by Print Electronics Is-
 76 rael [20] by mechanical drilling of 0.5 mm holes in FR4 plates,
 77 Cu-clad on one or two sides, followed by chemical etching of
 78 0.1 mm wide rims around each hole. In the double-stage detec-
 79 tors the THGEM had a hexagonal hole pattern with a pitch of
 80 1 mm; the SRWELL's square-shaped hole pattern had a pitch
 81 of 0.96 mm, with 0.86 mm wide "blind" strips above the grid
 82 lines (1.36 mm between the centers of the holes on the opposite
 83 sides of the strip). The resistive layers had a surface resistivity
 84 of 10-20 MΩ/square. The FR4 sheet serving as the base of the
 85 resistive anode was 0.1 mm thick. The grid patterned on the
 86 FR4 sheet had 0.1 mm wide copper lines, defining an array of
 87 8×8 squares, 1 cm² each, matching the 8×8 readout pad array
 88 patterned here on a 3.2 mm thick FR4 plate located below the
 89 anode.

90 For the data acquisition the new CERN-RD51-SRS electron-
 91 ics (Scalable Readout System [21]) was used, with the 64-
 92 pad array read by a single SRS analog 128-channel APV25
 93 chip [22]. External triggering and tracking were provid-

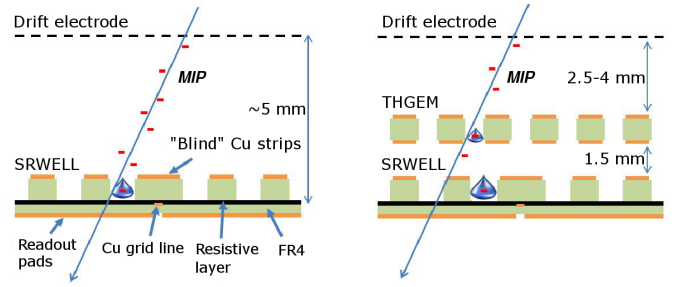


Figure 2: The two detector configurations investigated in this work. Left: single-stage SRWELL; Right: double-stage detector with a standard THGEM multiplier followed by an SRWELL.

94 ed using the RD51 tracker/telescope setup [23], comprising
 95 three 10×10 cm² scintillators in coincidence with three MI-
 96 CROMEGAS tracking units, each equipped with two APV25
 97 chips. The three tracker detectors and the investigated detector
 98 shared the same external trigger and front-end card (FEC), en-
 99 abling event-by-event matching and track reconstruction. This
 100 permitted measuring both the global average values of the de-
 101 tection efficiency and of the pad-multiplicity, as well as their
 102 local, position-dependent values (e.g. their variations at the pad
 103 boundaries). The low-noise electronics enabled operating the
 104 detectors at relatively modest gas gains of ~2000-3000.

105 The detector electrodes were biased individually through a
 106 CAEN SY2527 HV system. The voltage and the current of each
 107 HV channel were monitored and stored using the RD51 slow-
 108 control system [24], allowing for measuring the rate and the
 109 magnitude of occasional discharges (e.g. momentary voltage
 110 drops, accompanied by current pulses).

111 The detectors were investigated in a broad low-rate muon
 112 beam (10-20 Hz/cm²), and in narrow pion beams of ~1 cm²
 113 area. The pion rates were varied between ~0.5 kHz/cm² to
 114 ~70 kHz/cm², with the majority of the data taken at rates of
 115 up to a few kHz/cm².

116 Average and local values of the detector efficiency and pad-
 117 multiplicity were studied using selected tracker events. Pads
 118 were considered as activated if their signal was above a pad-
 119 specific threshold (individually set according to the noise level
 120 of each pad). The detector efficiency was defined as the fraction
 121 of tracks where a corresponding cluster of pads was found with
 122 its calculated center of gravity not more than 10 mm away from
 123 the track projection on the detector. These same tracks were
 124 used to calculate the average pad-multiplicity by counting the
 125 number of pads activated per event. For more details see [14].

3. Results

The studies on single-stage detectors included two configu-
 rations: one with a 0.4 mm thick SRWELL and a 5.5 mm drift
 gap, and the other with a 0.8 mm thick SRWELL and a 5 m-
 m drift gap. In a muon beam, the former reached 97% global
 efficiency at an average pad-multiplicity of 1.2, and the latter
 (0.8 mm thick SRWELL) displayed 98% global efficiency al-
 ready at 1.1 multiplicity. The measured Landau pulse-height

distributions were well above the noise level at gains of ~ 1500 – 171 2000. The discharge probabilities with muons were of the order 10^{-6} for both configurations. However, with pions both configurations displayed a gain drop by a factor of ~ 2 at the above operating conditions, with a ~ 5 – 10 fold increase in the discharge probability; this resulted in lower detection efficiencies with pions for both cases. In both detector configurations, the observed discharges could be divided into two distinct groups: (a) a vast majority of micro-discharges, involving small voltage drops (~ 10 – 15 V) with a typical recovery time of ~ 2 seconds; (b) a small fraction of discharges involving large voltage drops (~ 100 – 200 V) with longer recovery times (a few seconds, depending on the size of the voltage drop). The 0.8 mm thick SRWELL appeared to be more stable than the 0.4 mm thick detector, but this requires further study and more precise quantification.

The studies of the double-stage detectors were done with 0.4 mm thick THGEM and SRWELL electrodes. The transfer gap was kept at 1.5 mm and the drift gap was varied between 2.5 mm and 4 mm. The efficiencies recorded with muons were similar to those obtained with the single-stage detectors, albeit shifted to slightly higher multiplicities. For example, the 4 mm drift, double-stage detector reached 97% global efficiency at an average multiplicity of 1.15; the 3 mm drift, double-stage detector reached 94% efficiency at a multiplicity of 1.2. The discharge probabilities with muons were extremely low, of the order of 10^{-7} , for the 4 mm drift double-stage detector. Figure 3 shows the global efficiency versus the average pad-multiplicity for the 0.8 mm thick single-stage SRWELL detector and for the double-stage THGEM/SRWELL detector with 4 mm drift. Measurement details are provided in [14].

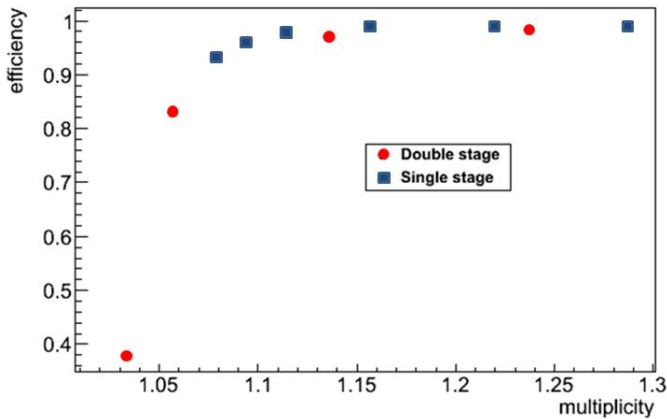


Figure 3: Global detection efficiency versus average pad-multiplicity for the 0.8 mm thick single-stage SRWELL detector with 5 mm drift gap, and for the double-stage THGEM/SRWELL detector with 4 mm drift gap and 1.5 mm transfer gap.

Unlike the results for the single-stage detectors, no gain drop was observed for the double-stage detectors when switching from muons to pions; Figure 4 compares the pulse-height distributions measured for both particle types with the double-stage detector having a 4 mm drift gap, under the same operation voltages. Although occasional discharges occurred with pions for

this detector, their probability, at rates of a few kHz/cm², was of the order of 10^{-6} , and the observed voltage drops (on the SRWELL top) were all minute and limited to ~ 3 V, with a recovery time of ~ 1 second (no large discharges were observed). The efficiency for pions was similar to the one obtained with muons ($\sim 95\%$ and above). A comparison between runs with and without these micro-discharges showed that their effect is negligible in terms of the detection efficiency. Moreover, the presence of micro-discharges had no effect on the data acquisition system, which operated stably even in high rate pion beams.

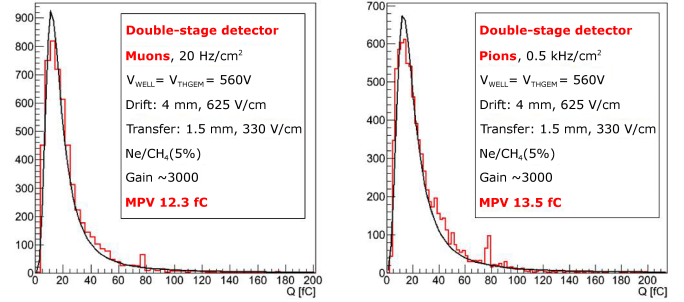


Figure 4: Pulse-height (Landau) distributions for muons (left) and pions (right) measured with the double-stage detector with 4 mm drift gap. The parameters and operation conditions are given in the figures. No gain drop was observed with pions in this double-stage configuration.

The ability to accurately match events between the tracker and the investigated detectors permitted studying the dependence of the local efficiency and of the pad-multiplicity on the track position relative to the pad boundary. The results are shown in Figure 5: essentially no drop in local efficiency occurred above the “blind” SRWELL strips in both configurations; the local increase in pad-multiplicity above the inter-pad boundary resulted from charge sharing between holes on the opposite sides of the copper strip (see Fig. 1).

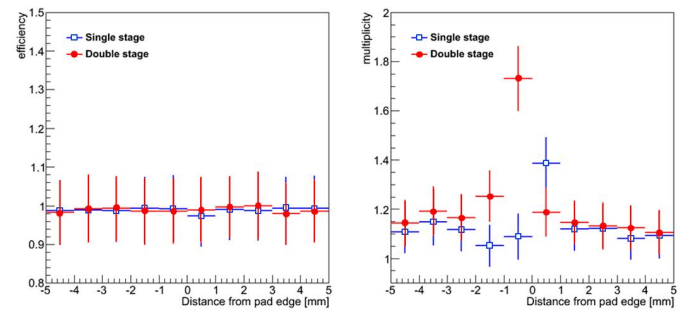


Figure 5: Local detection efficiency (left) and pad-multiplicity (right) as a function of the muon-hit distance from the pad boundary for the single-stage 0.8 mm SRWELL and the 4 mm drift double-stage detectors.

4. Summary and discussion

The beam tests described in this work were performed to investigate, for the first time, structures based on the Segmented Resistive WELL (SRWELL) concept. This new THGEM-variant has several key advantages which make it a promising

candidate for Digital Hadronic Calorimetry: (1) By removing the “standard” induction gap, it allows for a significant reduction in thickness - a critical feature in applications such as the SiD experiment; the total thickness of the detector configurations studied in this test was 5-6 mm (excluding the readout electronics); (2) The resistive anode effectively quenches occasional discharges, whose magnitudes, in the double-stage configuration, are limited to ~ 3 V with ~ 1 s recovery time - with effect on the detection efficiency or on the stability of the electronic readout system; (3) The copper grid underneath the resistive layer significantly reduces the cross-talk between neighboring pads, limiting the multiplicity to ~ 1.1 - 1.15 ; the higher value is mostly due to particles inducing avalanches on more than one hole; (4) The detection efficiency for muons is exceptionally high: 98% at a multiplicity of 1.1 with the single-stage 0.8 mm SRWELL, and 97% at a multiplicity of 1.15 with the 4 mm drift double-stage THGEM/SRWELL. Finally - the SRWELL, like the standard THGEM, is a robust electrode which is essentially immune to spark damage. It can be readily and economically produced for large areas, using industrial methods. The combination of the above properties make the SRWELL-based detectors highly competitive compared to the other technologies considered for the SiD-DHCAL.

Single-stage detectors are obviously advantageous in terms of cost when considering large-area applications such as the DHCAL. While their efficiency and multiplicity figures for muons are very convincing, the pion-induced gain drop in the single-stage SRWELL - not observed for the double-stage detectors - is intriguing, and should be clarified (and mitigated) in additional laboratory tests.

The detector thickness limitation imposed by the SiD experiment calls for the use of extremely thin front-end electronics (a requirement which is, at present, not met by the SRS system). Two alternative readout systems may be suitable for this application: SLAC’s KPIX board [25], already beam-tested with THGEM-based detectors [13], and the MICROROC chip developed by the LAL/Omega group and by LAPP [11], which was extensively tested with MICROMEGAS detectors. Investigations with THGEM-based detectors are already underway.

Optimization studies on SRWELL detectors (single- and double-stage), as well as work on larger detectors, are planned for the near future. One attractive option is the return to argon-based gas mixtures, implying 2-3 fold higher MIP-induced ionization electron numbers, though at the cost of higher operation potentials [26]. Modern low-noise electronics may allow for lower-gain operation, so this might be possible without problems.

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